Stimulated Raman spin coherence and spin-flip induced hole burning in charged GaAs quantum dots


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(Received 27 July 2007; published 12 March 2008)

High-resolution spectral hole burning (SHB) in coherent nondegenerate differential transmission spectroscopy discloses spin-trion dynamics in an ensemble of negatively charged quantum dots. In the Voigt geometry, stimulated Raman spin coherence gives rise to Stokes and anti-Stokes sidebands on top of the trion spectral hole. The prominent feature of an extremely narrow spike at zero detuning arises from spin population pulsation dynamics. These SHB features confirm coherent electron spin dynamics in charged dots, and the linewidths reveal spin spectral diffusion processes.

DOI: 10.1103/PhysRevB.77.115315

Single electron spin localized in semiconductor quantum dots (QDs) has attracted a great deal of interest due to its potential use in quantum applications.1 Experimental and theoretical efforts have been focused on controllable coherent spin dynamics and possible decoherence mechanisms.2–14 In this paper, we report spectral hole burning (SHB) in coherent differential transmission (DT) spectroscopy induced by spin-trion dynamics in an ensemble of negatively charged QDs. Spin-coherence induced SHB by stimulated Raman excitation and spin relaxation induced SHB due to the population pulsation dynamics are observed in the QD system. Features of SHB not only disclose important spin dynamics observed from transient spectroscopy and phase modulation techniques,11 but also provide information on spin spectral diffusion (SD) processes,14 which are not easily revealed by previous methods.

The interface fluctuation GaAs/Al0.3Ga0.7As QDs are molecular beam epitaxy grown with growth interrupts, and modulation Si doping in the barrier incorporates excess electrons.15 The pump E1(ω1) and probe E2(ω2) optical fields are derived from two frequency-stabilized and independently tunable cw lasers with a mutual coherence bandwidth of 20 neV, which is crucial for this experiment, as discussed below. The sample is kept inside a superconducting magnetic liquid helium flow cryostat, and the temperature is maintained at 4.5 K. The DT signal is homodyne detected with the probe field by a photodiode and extracted by a lock-in amplifier.

A nonlinear degenerate DT spectrum with ω1 = ω2 [Fig. 1(a)] shows the trion and exciton ensemble resonances. Their assignments are confirmed by both photoluminescence and transient quantum beat studies (data not shown). The trion binding energy (i.e., the separation between exciton and trion resonances) is measured to be 2.9 meV, in agreement with earlier reports of photoluminescence15 and transient spectroscopy.2 The ensemble trion inhomogeneous broadening width (~2.5 meV) can be estimated from the broad Gaussian profile of the trion resonance.

To study SHB with nondegenerate DT spectroscopy, the pump beam is fixed near the trion ensemble peak and the probe beam is detuned (Δ = ω2 − ω1). A narrow spectral structure appears, exhibiting a double-Lorentzian-like shape, i.e., a narrower Lorentzian peak on top of a broader one [Fig. 1(c)]. As established below, the narrower Lorentzian peak is due to the trion population relaxation dynamics, and its linewidth (~12 µeV) gives the trion population relaxation rate.

FIG. 1. (Color online) (a) Coherent degenerate DT spectrum of both trion and exciton ensemble resonances. The arrow shows the fixed pumping position for (c). (b) Energy level diagram of charged QDs in the absence of magnetic field, where solid (empty) circles indicate electrons (holes) and arrows show the corresponding spin orientations. The single electron ground state has either spin down or spin up. The trion state consists of a two-electron singlet state and one heavy hole with spin down (|r−⟩) or spin up (|r+⟩). σ+σ−(σ+) light couples |z−⟩ (|z+⟩) to |r−⟩ (|r+⟩). (c) A clear view of the trion spectral hole burning profile, shown as a double Lorentzian, where the probe is detuned relative to pump position (1621.9 meV).
\[ i\hbar \dot{\rho}(\epsilon) = [H, \rho(\epsilon)] + \frac{\partial \rho(\epsilon)}{\partial t}_{\text{relax}} + \frac{\partial \rho(\epsilon)}{\partial t}_{\text{SD}}, \tag{1} \]

where \( \rho(\epsilon) \) is the density matrix of each ensemble member characterized by its resonance energy \( \epsilon \). \( H \) is the total Hamiltonian including the interaction with the coherent optical fields. The second term on the right hand side is a generalized relaxation term that describes population decay and pure dephasing. The last term is due to various SD processes. Equation (1) without the SD term reduces to the standard OBE.

Considering the optical selection rules in the absence of a magnetic field, the transition from spin ground state \( |z^-\rangle \) to trion state \( |l^+\rangle \) forbidden as a dark transition. Therefore, with circularly polarized pump and probe fields, the charged QD energy levels reduce to a two-level system, as shown in the enclosed box in Fig. 1(b). The relevant SD process here is the interdot transfer of trion population,

\[ \frac{\partial \rho_{J}(\epsilon)}{\partial t}_{\text{SD}} = -\Gamma_{J}^{SD}(\epsilon)\rho_{J}(\epsilon) + \int W_{J}(\epsilon, \epsilon')\rho_{J}(\epsilon')d\epsilon', \]

where \( W_{J}(\epsilon, \epsilon') \) is the spectral redistribution kernel representing the rate for trion population to migrate from a dot with resonant energy \( \epsilon' \) to dots with energy \( \epsilon \). While this interdot transfer conserves the ensemble trion population, the dipole coherence is lost. \( \Gamma_{J}^{SD}(\epsilon) = \int W_{J}(\epsilon, \epsilon')d\epsilon' \) is the overall SD rate. For an optically thin sample, the DT spectrum is given by the imaginary part of the third order induced optical polarization integrated over the inhomogeneous distribution.

\[ P_{\text{NL}}^{(3)} = -\frac{iN\mu|E|^2E_{z}^{2}}{2\hbar^2\Lambda_{i}} \left[ \frac{1}{\Gamma_{i}^{SD} + i\Delta} \right] \times \left[ \frac{1}{\Lambda_{i}} + \frac{\sqrt{\pi}}{2(\gamma_{i} + \Gamma_{i}^{SD}) + i\Delta} \right] \left[ \frac{\Gamma_{i}^{SD} - \sqrt{\pi}}{\gamma_{i} + \Gamma_{i}^{SD} + i\Delta} \right] + \frac{1}{\Gamma_{i}^{SD} + i\Delta} \right] \]

where \( \Gamma_{i}(\gamma_{i}) \) is the trion population (coherence) decay rate, \( N \) is the total number of excited charged QDs, \( \Lambda_{i} \) is the ensemble inhomogeneous broadening width of the trion resonance, and \( \mu \) is the optical dipole of a single dot. We have assumed a redistribution kernel \( W_{J}(\epsilon, \epsilon') = \Gamma_{J}^{SD} \exp\left[-(\epsilon - \epsilon')^{2}/(\Lambda_{i})^{2}\right] / (\sqrt{\pi}\Lambda_{i}) \), where \( \bar{\epsilon} \) is the ensemble averaged resonance energy.

Equation (3) holds when the approximation of plasma dispersion function is taken because \( \Gamma_{i}, \gamma_{i}, \Gamma_{i}^{SD}, \Delta \ll \Lambda_{i} \) and the pump frequency is fixed near the center of the ensemble trion spectrum.

The SD process significantly changes the trion SHB line shape and linewidth, which can be discussed in three regimes, as schematically shown in Fig. 2. \( \Gamma_{i} = 2\gamma_{i} \) is assumed only for the theoretical calculation in Fig. 2 to simplify discussions since pure dephasing has been found negligible for trion states at 4.5 K. When \( \Gamma_{i}^{SD} < \Gamma_{i} \), the trion SHB in the DT spectrum reduces to the standard Lorentzian squared with width of \( \sim \Gamma_{i} \). Second, when \( \Gamma_{i}^{SD} > \Gamma_{i} \), the trion SHB line shape is \( (2\gamma_{i}^{2} + \Delta^{2})/[(\Gamma_{i}^{SD} - \Delta^{2})^{2}] \), where the linewidth is considerably broadened by the SD process. Finally, when \( \Gamma_{i}^{SD} \approx \Gamma_{i} \), the line shape changes to a Lorentzian-like profile with larger full width at half maximum (FWHM) of \( 2\gamma_{i}^{2} \) and smaller FWHM of \( 2\Delta^{2} \). The physical origin of the narrower Lorentzian is due to the population pulsation effect.

The observed line shape shown in Fig. 1(c), which is consistent with the physical model [Eq. (3)] assuming \( \hbar \gamma_{i} = 7 \mu eV, \ h\gamma_{i} = 8 \mu eV, \ h\gamma_{i} = 46 \mu eV, \ h\gamma_{i} = 1000 \mu eV \) plus a slow linear slope. The values of \( \Gamma_{i} \) and \( \gamma_{i} \) agree with earlier reports, and the value of \( \Lambda_{i} \) is in agreement with the degenerate DT trion ensemble resonance in Fig. 1(a). The relatively big \( \Gamma_{i}^{SD} \) implies that the trion SD process plays an important role.

With a magnetic field applied perpendicular to the QD growth direction (i.e., Voigt geometry), the spin eigenstates \( |x \pm, y \rangle \) along the field direction are split in energy due to the nonzero electron in-plane \( g \) factor \( (g_{l}^{e}) \) while the trion states are unaffected as the heavy-hole in-plane \( g \) factor is negligible. The new eigenstates and optical selection rules are shown in Fig. 3(a). With circularly polarized pump and probe fields, the relevant states form a three-level \( \Lambda \) system enclosed in the dashed box. Utilizing a similar \( \Lambda \) system, optical pumping leading to spin cooling has been reported recently in different structural QDs. However, no similar optical pumping effect was observed in our system, where the strong spectral diffusion on two spin ground states discussed below plays an important role.

The SHB profile changes dramatically in this Voigt geom-
FIG. 3. (Color online) (a) Energy level scheme of charged dot in Voigt geometry \((B_z≠0)\). The new ground states \(|x±\rangle\) denote electron spin population along the \(x\) axis separated by Zeeman splitting. The new selection rules are labeled by solid (dashed) lines with \(σ^z(σ^z)\) light. (b) The newly emerged center sharp spike and two symmetric sidebands at \(B_z=2.2\ T\), where the fixed pump is indicated by the arrow. The two Stokes and anti-Stokes spin coherence induced sidebands are highlighted. The inset shows the zoomed in sharp central spike at \(Δ=0\). A similar figure was first presented in Ref. 11, but with a different objective.

FIG. 4. (Color online) Comparison of theoretical calculations and experimental results in Voigt geometry. Theory: The DT spectrum of a three-level \(Λ\) system at different magnetic fields, where \(Λ_3=1000Γ_s\), \(γ_s=Γ_s/2\), and \(Γ_s^{SD}=0.01Γ_s\). (a) \(B=0\), so that spin inhomogeneous broadening is due to nuclear field \(Λ_3−Λ_s^0\). The central spike and the sidebands merge in the limit of vanishing \(B\). (b) \(B=2B_0\), where \(μ_Bg_sB_0=3hΓ_s\). Spin inhomogeneous broadening is dominated by the inhomogeneity of the \(g\) factor: \(Λ_3=μ_BΔg'_sB_s/h\). (c) \(B=2B_0\). Experiment: SHB line shapes at different magnetic fields, where the arrow indicates the fixed pump position (1622 meV) and the insets show the zoomed view of central sharp peaks.

is the redistribution kernel and \(Γ_s^{SD}(ε_s,ε_s)\)
\(=\int W_s(ε_s,ε_s,ε_s,ε_s) dε_s' dε_s''\) is the spin SD rate. The qualitative feature of the redistribution kernel function depends critically on the SD mechanism. Local nuclear field fluctuation induced SD only affects the spin Zeeman splitting: \(W_s(ε_s,ε_s)\). We note that the inhomogeneous broadening of \(ε_s\) induced by the nuclear field is given by \(\Lambda_3≈−0.1\ \mu\) eV \(≪\gamma_s\). We also note that two quantum dots are equally excited if the difference in their resonance frequency for spin to trion transition is much smaller than the trion broadening \(γ_t\). Thus, \(ρ_+^{(0)}(ε_s,ε_s)\) \(=\rho_{−}^{(0)}(ε_s,ε_s)\) if \((ε_s−ε_s')≤(ε_s−ε_s)/2\). It can then be shown that the two terms on the right hand side of Eq. (5) cancel each other, and hence this mechanism has a negligible effect on the linewidth of the ultranarrow central spike associated with the spin population pulsation dynamics. On the other hand, if the SD process is due to the interdot transfer of nonequilibrium spin population, it is more reasonable to assume a redistribution kernel \(W_s(ε_s,ε_s)\) \(=\int W_s(ε_s,ε_s,ε_s,ε_s) dε_s' dε_s''\) if \((ε_s−ε_s')±(ε_s−ε_s)/2\). In the vicinity of zero detuning \(Δ=0\), the DT signal is determined by

\[
E_{NL} = \frac{N\gamma_s|\mu_0|^2|E|^2E^*}{8\hbar^3\Lambda_3(2γ_s+Γ_s^{SD})Δ^2 + (2Γ_s+Γ_s^{SD})^2},
\]

shown as the sharp central spike with linewidth of \(2Γ_s^{SD}\) in Fig. 4 (Theory). The experimental data at various magnetic fields are shown in Fig. 4 (Experiment). and the measured linewidth of the sharp central spike is plotted in Fig. 5(b) from which we extract \(hΓ_s^{SD}≈0.2\ \mu\) eV.

In addition to the sharp central spike, the newly emerged SHB features in Voigt geometry also include two symmetric sidebands as highlighted by the ellipse regions in Fig. 3(b).
Stokes and anti-Stokes sidebands appear respectively at \( g/e = \frac{\Delta B}{h} = \frac{\gamma_s}{h} \) rate. The ensemble averaged sideband line shapes are a coherence, it can be inferred from Eq. (9) the sideband feature is associated with the Raman spin conservation. Moreover, a spin SD process is observed in contributions to the SHB line shape and has been theoretically identified to result as consequence of spin population pulsation dynamics and stimulated Raman spin coherence, respectively. In summary, we have shown in this paper that SD from the trion state complicates the trion SHB profile and makes an important contribution to the double-Lorentzian-like line shape. In the Voigt geometry, we have found a complex line shape arising from spin dynamics in the spectral hole burning: a narrow central spike and two symmetric Stokes and anti-Stokes sidebands. They have been theoretically identified to result as consequence of spin population pulsation dynamics and stimulated Raman spin coherence, respectively. Moreover, a spin SD process is observed in contributions to the SHB line shape and has been theoretically identified as an interdot transfer of the nonequilibrium spin population. Possible mechanisms include the interdot spin flip-flop interactions of the electrons or spin conserved electron tunneling to adjacent neutral dots. Their quantitative estimates cannot be determined without a detailed calculation, which is beyond the scope of this paper. Nevertheless, the existence of these decoherence mechanisms revealed by our experiments will have important impacts on the efforts toward spin based quantum applications in these kinds of QDs.

ACKNOWLEDGMENTS

This work was supported in part by the U.S. ARO, NSA/LPS, ARDA, AFSOR, ONR, and FOCUS-NSF.
21 We have assumed the strong redistribution model for interdot trion and spin spectral diffusion where the kernel only depends on the final state density (Ref. 16).