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Laser-induced surface acoustic waves: An alternative method to nanoindentation for the mechanical characterization of porous nanostructured thin film electrode media



MATERIALS

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ABSTRACT

The mechanical characterization of electrode materials in thin film lithium ion batteries is currently a sparse area. However, mechanical studies could offer valuable insight since the performance and breakdown of active materials is electromechanically coupled. In this paper, a porous nanostructured V₂O₅ cathode thin film with demonstrated high electrochemical performance was investigated by a laser-induced surface acoustic wave technique (LiSAW) that mitigates some of the challenges associated with the popular nanoindentation technique. The intent was to explore the capability of LiSAW in measuring the elastic modulus of the nanostructured film such that a reliable methodology could be produced to mechanically characterize challenging electrode materials. LiSAW measured a modulus of 53 ± 4 GPa for the porous V₂O₅ film and had no problems coping with the 40 nm roughness and delicate structure. On the other hand, nanoindentation produced a modulus of 50 ± 10 GPa, which is comparable to LiSAW, but with considerably higher uncertainty from roughness. For porous nanostructured electrodes, and other challenging films, that are too soft, thin, or delicate for traditional nanoindentation measurements, LiSAW is a potentially excellent alternative. LiSAW testing on many other electrode materials would be instrumental in developing a better understanding between the mechanical and electrochemical properties of thin film battery materials.

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1. Introduction

Lithium ion batteries are an integral part of today's energy storage landscape, but the constant miniaturization of electronic devices is currently challenging the technology to fit into reduced form factors. While smaller electronics have allowed for great advancement in many areas such as portable consumer electronics and biologically implantable devices (Holmes, 2001), their energy requirements can often remain similar, if not greater than,

http://dx.doi.org/10.1016/j.mechmat.2015.10.005 0167-6636/© 2015 Elsevier Ltd. All rights reserved. before (Whittingham, 2012). Thin film Li-ion batteries are capable of meeting the limited size requirements, but alternatives to the traditional graphite and lithium metal oxide electrode materials are necessary to maintain capacity and performance with diminished battery mass (Ji et al., 2011; Ellis et al., 2010). Nanostructured materials are of interest in this field because they can exhibit very high surface energies which enhance charge transfer kinetics and ion storage capacity (Liu and Cao, 2010; Song et al., 2011; Zhang et al., 2013). Additionally, nanostructures, especially when coupled with porosity, can be more tolerant to the large physical deformations associated with Li-ion insertion and removal thus reducing the mechanical



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breakdown of the electrode and improving the battery life cycle (Vu et al., 2012; Hayner et al., 2012).

It is well known that the effects of the electrode volume change during each charge cycle, producing modest strains of 2% in LiCoO₂ cathodes (Reimers and Dahn, 1992) and 7% in LiFePO₄ cathodes (Meethong et al., 2007), contributes greatly to the loss of charge capacity over time (Cabana et al., 2010). Yet for such a mechanically rich process, there has been relatively little investigation correlating the breakdown to the mechanical properties of the electrode. While many studies have targeted and identified the main degradation mechanism, the proliferation of micro-cracks (Ebner et al., 2013), comprehensive mechanical design criteria to mitigate such issues in new electrode materials have yet to be determined. With many other electrochemical considerations to contend with, it is easy to see why mechanical aspects are often overlooked, but with links between mechanics and rate capability (Meethong et al., 2007) and the emergence of higher capacity materials with larger expected volume changes, nanomechanical characterization may become critical in understanding optimal electrode design.

The prevalence of nanoindentation in thin film mechanical analysis is especially valuable to the characterization of thin film electrodes. Indentation tests can quickly determine the elastic and plastic properties of a material with the elastic modulus being especially useful here. While electrode degradation is a complex problem with both electrochemical and mechanical considerations, the modulus can breathe insight into the flexibility and reliability of a material and thus aid in understanding how it may cope with lithiation induced strain. Several recent studies have successfully applied nanoindentation to electrode materials. Qu et al. reported on the elastic modulus, density, and fracture toughness of individual LiCoO₂ grains and noted that the experimentally determined modulus of 174 GPa was much lower than the previously predicted value of 315-516 GPa (Ou et al., 2012). This suggests a more compliant structure than once thought. Ramdon and Bhushan reported on the elastic modulus, hardness, and wear properties of LiFePO₄ cathodes before and after multiple charge cycles (Ramdon and Bhushan, 2014). They found that the modulus did not change significantly through cycling, but that the hardness increased by almost 100% giving rise to a durable, yet brittle structure. Zhu et al. performed indentations on RuO₂ anodes and found that the elastic modulus decreases by an order of magnitude after 50 cycles (Zhu et al., 2013). With such a limited quantity of studies, it is difficult to correlate the elastic modulus, or any other mechanical property, to battery performance over time therefore this area demands further investigation.

The lack of nanomechanical studies could possibly be attributed to the difficulty of nanoindentation on thin film electrode materials. For consistent results that are free of indentation substrate effects, it is necessary to have films with low surface roughness and a thickness many times that of the penetration depth. This can be especially problematic to nanostructured or porous films since their engineered roughness may be too delicate for mechanical polishing and their thicknesses could be very small. While substrate effects from low thicknesses can be overcome by utilizing an alternative indentation analysis scheme (Li and Vlassak, 2009), albeit with significant computational intensity, defeating roughness through deep indents may not always be viable. In this scenario, it may be more efficient to use a laser-induced surface acoustic wave (LiSAW) technique to determine the elastic modulus.

The LiSAW technique is valuable to thin and porous films because it requires only elastic deformations, it features non-contact measurements, and the mechanical property extraction naturally considers and isolates substrate coupling (Hess, 1996; Schneider and Tucker, 1996; Xiao et al., 2011). In the experiment, an Nd:YAG laser is used to excite Rayleigh waves that subsequently propagate across the layered system. Due to elastic mismatch between the film and substrate, frequency dispersion of the wave occurs. Since the relationship between mechanical properties and dispersion is well documented in wave mechanics (Farnell and Adler, 1972), LiSAW can seek unknown parameters of the film or substrate by experimentally quantifying the dispersion. In this manner, if enough nonlinearity exists in the dispersion curve, the film thicknesses, Poisson's ratios, densities, and elastic moduli of the materials involved can be found. Typically, many of these parameters are known at the onset, thus only one or two parameters may require fitting. LiSAW can deliver these results on the same time-scale as typical indentation tests and with no substrate influence. Due to the micron scale of SAW wavelengths, the technique can work well on moderately rough films (Cote et al., 2009) and, depending on SAW detection scheme, with little or no surface preparation. While the LiSAW technique requires an area millimeters in scale for adequate wave dispersion with existing detection schemes, the creation of uniform blanket films of electronic materials is typically not a major challenge. Rather, it is the thickness of such films that creates the greatest characterization challenge and LiSAW directly addresses and excels at these geometries based on the inherently coupled film-substrate analysis procedure.

The limiting factor of many current Li-ion batteries lies in the low energy density of the cathode. Common cathode materials based on layered lithium metal oxides such as LiCoO₂ and LiFePO₄ have capacities of just 140 mA h g^{-1} and 170 mA h g⁻¹ respectively (Nazri and Pistoia, 2009). Significant research has been completed on the oxides of other first row transition metals in order to achieve higher capacities. From these studies, vanadium oxides have emerged with great potential due to their wide range of available oxidation states, high reversibility, and favorable layered structures, all of which are conducive to the electrochemical lithiation process (Ding et al., 2009). Specifically, V_2O_5 with a theoretical capacity of 450 mA h g⁻¹ (Liu et al., 2011a), has been extensively studied and remains a leading candidate amongst alternative cathode materials. While V₂O₅ has been known for inadequate structural stability and slow electrochemical kinetics, nanostructured V₂O₅ forms have recently been shown to overcome these challenges by demonstrating excellent capacity, rate capability, and cyclability (Liu et al., 2011b; Pomerantseva et al., 2012; Augustyn and Dunn, 2010).

In this paper, the mechanical properties of a porous nanostructured V₂O₅ cathode film are investigated by nanoindentation and LiSAW. This particular film is especially attractive because it is synthesized by a simple and cost effective procedure via electrodeposition and exhibits excellent electrochemical properties. From a mechanical standpoint, the film's roughness and delicate nature, both properties common to nanostructured electrode materials and for which mechanical polishing is not viable, makes it an interesting specimen for testing the limits of traditional nanoindentation. The goal of this study is to perform LiSAW and discover whether the technique is reliable for such films. The results will then be compared to findings from nanoindentation with the intent of identifying the most efficient means for mechanical characterization of such media. It is hoped that the methodologies identified from this research will guide and fuel additional investigations in order to develop a further understanding of the relationship between electrode performance and its mechanical properties.

2. Experimental method

2.1. Electrodeposited V₂O₅ film synthesis

For mechanical characterization, electrodeposited V₂O₅ films were prepared on single crystal silicon wafers with (100) orientation by the procedure reported in literature (Liu et al., 2011a). In brief, a vanadium pentoxide solution was formed by combining commercial V₂O₅ powder (99.8% Alfa-AESAR) with H₂O₂ (30 wt.% in water, Sigma-Aldrich) and de-ionized water. The solution was stirred for 15 min and then sonicated for 15 min in water at constant temperature. Afterwards, the V₂O₅ concentration was diluted from $C_V = 0.3$ M to $C_V = 0.06$ M and the resulting solution was sonicated for 1 h. A final dilution, this time to $C_V = 0.0075$ M, completed the V₂O₅ solution and prepared it for use in electrodeposition.

Prior to electrodeposition, the Si substrates were sputter coated with less than 5 nm thick gold palladium to enhance conductivity. They were then lowered into the V₂O₅ solution parallel to a platinum plate counter electrode. Deposition was carried out at -2.4 V for 1 h with a separation distance between the electrodes of ~ 15 mm. Upon completion, the films were dried in air at 70 °C for 1 h and then annealed in air at 450 °C for 3 h. To obtain free standing films for density measurements, the procedure was carried out on nickel foil substrates (99.9% Sigma-Aldrich) instead of silicon. After electrodeposition, the undried gel films were carefully detached from the foil and bridged across a \sim 20 mm gap on wax paper. The films were dried and cut free from the wax paper to isolate their spanning areas prior to annealing. Drying and annealing conditions remained unchanged from before. All film thicknesses were measured by profilometry (Dektak, Tucson, AZ).

2.2. Electrochemical analysis of V₂O₅ films

For electrochemical analysis, V₂O₅ thin film electrodes were deposited using the same procedure described above

but with a deposition time of 30 min on FTO glass and annealing temperature of 500 °C. The slightly higher temperature was chosen to assure the complete removal of the crystalline water in the V₂O₅ interlayers since nonaqueous solvent is used in the electrochemical testing. However, the 50 °C difference in annealing temperature is not expected to produce any discernable difference in the microstructure and mechanical properties of the film. The electrodes were tested using a standard three-electrode configuration with 1 M LiClO₄ (98%, Sigma Aldrich) in propylene carbonate (99.7%, Sigma Aldrich) as electrolyte, a Pt plate as counter electrode, and Ag/Ag⁺ as reference electrode (CH Instruments). Both cyclic voltammetry (CV) and chronopotentiometry (CP) tests were performed using an electrochemical analyzer (CH Instruments, Model 605B). CV was conducted in the voltage range of 0.6 to -1.8 V (vs. Ag/Ag⁺) at a scan rate of 5.0 mV s⁻¹. CP was also performed in the voltage range of 0.6 to -1.8 V (vs. Ag/Ag⁺) assuming a 1 C current density of 298 mA g^{-1} corresponding to the total insertion of two lithium ions per [V₂O₅] formula unit.

2.3. Mechanical characterization by nanoindentation

Nanoindentation measurements were performed by a Hysitron UBi1 nanomechanical test instrument (Hysitron, Minneapolis, MN) with a diamond Berkovich tip (142.3° total included angle). All indentations followed a trapezoidal loading curve featuring a 10 s ramp loading segment, a 5 s hold at maximum load, and a 10 s unloading period. To decrease the likelihood of substrate influence, maximum indentation depths were contained to 10% of the film thickness when possible. Mechanical properties were derived from indentation loading curves by the common Oliver and Pharr method (Oliver and Pharr, 1992).

2.4. LiSAW experimental setup

LiSAW measurements were performed by the apparatus shown in Fig. 1. The setup is highlighted by a 5 ns pulse width Nd:YAG laser (New Wave Research, Fremont, CA) operating at 1064 nm for SAW generation and a continuous wave argon-ion laser (Coherent, Santa Clara, CA) operating at 514.5 nm for SAW detection. For planar SAW generation, the circular Nd:YAG beam was focused by a cylindrical lens to a line source on the specimen surface with a width of ${\sim}8\,\mu\text{m}$. During SAW detection, out-of-plane surface displacements were measured by the argon-ion laser in a Michelson interferometer configuration. Displacement for this type of interferometer is obtained through laser interference as the light is split into two paths by the beam splitter, one to the sample and one to a reference mirror, then later recombined. To maximize frequency detection, the beam was expanded prior to focusing by two identical short-focal distance lenses to achieve a circular spot size of <10 µm on the sample and reference mirror. The interference signal was detected by a 1.2 GHz photo detector (EOT, Traverse City, MI) with a 300 ps rise/fall time and then digitized by a 5 GHz oscilloscope sampling at a rate of 20 gigasamples/s (LeCroy, Chestnut Ridge, NY).



Fig. 1. LiSAW experimental setup. Abbreviations: DG: delay generator, O-scope: oscilloscope, PD: photo detector, CXL: convex lens, CCL: concave lens, BE: beam expander, BS: beam splitter, M-PZ: mirror mounted on piezoelectric stage, CYL: cylindrical lens.



Fig. 2. Schematic of synchronization scheme that allows for mid-fringe laser firing for optimal displacement sensitivity.

In a Michelson interferometer, the voltage signal recorded by the oscilloscope will oscillate as the sample moves in time. The signal is directly related to sample displacement such that a sample movement of $\lambda/4$, where λ is the laser wavelength, represents half of a fringe (i.e. maximum destructive interference to maximum constructive interference). Since SAW displacements were much lower than $\lambda/4$ for our setup, the detection was considered subfringe. In this regime, mid-fringe measurements were necessary since in that region, the displacement sensitivity is considerably higher than in other regions of the fringe. This is because of the sinusoidal nature of fringe intensity with respect to surface displacement. Near the peaks and troughs of fringes, sample displacement produces less change in light intensity since the slope is nearer to zero. Conversely, between the peak and trough (mid-fringe), the slope is highest thus generating the largest change in light intensity per unit of displacement.

SAW events were synchronized to the middle linear fringe region with a technique similar to that reported by

the authors previously (Chow et al., 2014). In brief, a piezoelectric stage was mounted on the reference mirror and to set to oscillate. This produced an interference pattern with consistent and repeatable temporal wavelengths (Fig. 2) that could be tuned to the timescale of the generation laser firing operation (several hundred microseconds) by adjusting the piezoelectric vibration frequency. An oscilloscope monitoring signal rise and fall times could then detect selected wavelengths and use them to trigger and time the laser output to the desired linear fringe region. This synchronization technique relies on the ability to oscillate the mirror fast enough to overcome background noise, yet slow enough so as to not impede on the SAW timescale. Typically, a frequency of 100–200 Hz with a mirror movement of several microns was sufficient.

2.5. SAW Signal processing and mechanical property extraction

The detected SAW signals were inherently noisy with contamination from background vibrations and high frequency electronics. To suppress this noise, multiple signals were averaged together and a Fourier filter with a passing band of 15–300 MHz was used. To isolate the SAW event, a Tukey window function was employed with β = 0.25. Since SAW dispersion presents as a physical waveform change during propagation, two SAW signals, with known spacing from one another, are required for dispersion analysis. After obtaining the signals on each sample, dispersion was calculated through Eq. (1):

$$c(f) = \frac{2\pi f(d_2 - d_1)}{\Phi_2(f) - \Phi_{11}(f)} \tag{1}$$

where *c* is phase velocity, *f* is frequency, $d_2 - d_1$ is the spacing between detection sites, and Φ is the phase spectrum for each signal. This produced the characteristic dispersion plot of phase velocity versus frequency.

To extract mechanical properties, the experimental dispersion curves were fit with theoretically generated curves where the elastic modulus or density was the fitting parameter. Since only linear dispersion was seen, only one parameter could be fit. The derivation of theoretical dispersion curves with assumed substrate and film mechanical properties followed the same procedure as detailed in literature (Farnell and Adler, 1972). In brief, particle displacement was modeled with a generalized wave equation (Eq. (2)) and coupled to material properties through the governing equation of motion (Eq. (3)).

$$u_j = \alpha_j \exp(ikbx_3)\exp[ik(x_1 - \nu t)]$$
⁽²⁾

$$\rho \frac{\partial^2 u_j}{\partial t^2} = c_{ijkl} \frac{\partial^2 u_k}{\partial x_i \partial x_l} \quad (i, j, k, l = 1, 2, 3)$$
(3)

In Eqs. (2) and (3), u_j is displacement, α_j is relative amplitude, k is the wave vector, b is a complex parameter that governs wave decay, v is wave velocity, t is time, and c_{ijkl} is the stiffness tensor. By inserting Eq. (2) into Eq. (3), a system of 3 equations results for each material that contains an eigenvalue problem that can be solved for b with assumed wave velocity and mechanical constants. With *b*'s known for all materials involved, boundary conditions are considered that are comprised of stress and displacement continuity at the interface and zero stress at the free surface. The boundary conditions form a set of equations that can be solved together to seek the wave vector, k, which houses frequency information. The analysis is repeated for many phase velocities to form a theoretical dispersion plot of phase velocity versus frequency for assumed mechanical properties. In the cases where a secondary reflective layer was used for facilitating optical detection, the derivation followed a similar procedure (Xiao and You, 2006) to account for the second top film. The latter derivation applies to this work.

3. Results and discussion

3.1. Structure confirmation and electrochemical performance

Two samples were synthesized for LiSAW testing with each having dimensions of $\sim 25 \times 25$ mm, while the sample synthesized for electrochemical analysis had dimensions of $\sim 10 \times 10$ mm. The electrodeposited films were characterized by X-ray diffraction to confirm the presence of V₂O₅. The resulting spectrum, shown in Fig. 3, clearly displays the prominent (001) and (101) peaks of orthorhombic V_2O_5 thus confirming the oxide structure. Fig. 4 shows the rough nanostructured surface as imaged by a scanning electron microscope. The film thicknesses for mechanical characterization were measured to be 1.7 µm and the film thickness for the electrochemical analysis is \sim 800 nm. The two sets of films are structurally identical. It was the intent of the authors to create thicker films for the mechanical characterization in order to isolate the challenge of nanostructured roughness from the substrate effect during latter indentation tests.

Nitrogen adsorption isotherms were collected in order to evaluate the surface area and pore characteristics. The Barrett–Joyner–Halenda (BJH) pore size distributions



Fig. 3. X-ray diffraction spectrum of the nanostructured V_2O_5 film compared to a reference spectrum (National Bureau of Standards US, 1959).



Fig. 4. Image of the $V_2 O_5$ surface taken by a scanning electron microscope.

obtained suggest that the sample contains broadly distributed pores, the majority of which are smaller than 15 nm, with an average pore size of 3.19 nm; the BJH pore volume was 0.054 cm³ g⁻¹ while the Brunauer–Emmett–T eller (BET) derived specific surface area was 32.58 m² g⁻¹. The nanoscale pore size does not pose challenges to characterization by LiSAW and nanoindentation since both techniques involve microscale interactions, such as the physical wavelengths in LiSAW and tip radii in nanoindentation, which are both multiple orders of magnitude larger than the pore diameter.

To confirm the electrochemical performance, CV was performed to elucidate the redox processes at play for the V₂O₅ thin film electrodes, the first and eighth cycles of which are shown in Fig. 5(a). There are at least three clearly observable reduction and oxidation peaks for the initial CV scan, but these peaks become attenuated and broadened with further cycling; such findings are typically attributed to a loss of electro-active material crystallinity due to the lithium insertion/extraction process (Cocciantelli et al., 1991). Regardless, the reduction peaks that accompany lithiation at -0.20, -0.40, and -1.65 V vs. Ag/Ag⁺ can be ascribed to several well documented V_2O_5 phase transitions, namely α/ϵ , ϵ/δ , and δ/γ , respectively. Conversely, the oxidation peaks at -0.7, -0.18, and -0.008 V vs. Ag/Ag⁺ correspond to the γ/δ , δ/ε , and ε/α phase transitions, respectively (Braithwaite et al., 1999; Delmas et al., 1991; Galy, 1992). Comparison of these two cyclic voltammograms also reveals a slight drop in the overall current with cycling that can be ascribed to solid electrolyte interphase (SEI) formation. The formation of this organic and inorganic dual layered product, formed through electrolyte decomposition, effectively sequesters and intercepts lithium, leading to the subsequent drop in current (Ratnakumar et al., 2001). The potential difference between the oxidation and reduction peaks of each redox couple, which is a relative measure of the reversibility of the redox reaction, also diminishes with cycling. These considerations ultimately prove that the electrode is becoming stable in the later stages of cycling.



Fig. 5. Electrochemical performance of the 500 °C annealed V_2O_5 thin film electrodes: (a) cyclic voltammogram at a scan rate of 5 mV s⁻¹ for the initial and eighth cycles, (b) rate capability and cyclic stability at various current densities.

Fig. 5(b) shows the discharge capacity at a current density of 200 mA g^{-1} (0.67 C) for the first 20 cycles, followed by 120 cycles of discharge/charge tests at various current densities. The initial discharge capacity is 410 mA h g^{-1} , which surpasses the theoretical capacity corresponding to the insertion of two Li⁺ per V₂O₅ formula unit but is in actuality an inflated value because of SEI layer formation. Following this, the second discharge capacity dropped to 352 mA h g^{-1} , and then stayed at 320 mA h g^{-1} up to the 20th cycle. The coulombic efficiency for the initial and twentieth cycles is 83.2% and 98.9%, respectively. The initial discharge capacity at current densities of 0.67, 1.5, 5, 15, and 35 C were 279, 220, 164, and 126 mA h g⁻¹, respectively. After more than 130 cycles tested at different current densities, the current rate was returned to 1.5 C and the discharge capacity was still as high as $256 \text{ mA} \text{ h g}^{-1}$. The specific energy densities were calculated to be 700 W h kg¹ for the twenty-first cycle and 670 W h kg⁻¹ for the 150th cycle when discharged at 450 mA g^{-1} (1.5 C).

The high electrochemical performance demonstrated here could be ascribed to the nanometric scale of the V_2O_5 thin film electrodes which can provide a shorter diffusion path for Li-ion intercalation/deintercalation, and a higher surface area that offers more accessible intercalation sites which favors electrolyte penetration and interfacial reactions. As evidenced by the LiSAW results later, the externally applied electrical field during deposition assists adhering the film to the conductive substrate which facilitates Li-ion diffusion and enables excellent energy and power density; the improvement in mechanical integrity is especially apparent at high current densities.

3.2. Mechanical properties by LiSAW

When performing the electrodeposition, low resistivity silicon wafers were used for substrates. Heavy doping, with impurity concentrations on the order of 10^{19} cm⁻³, was necessary to achieve resistivities <0.002 Ω -cm and this has been shown to modify the elastic constants of silicon (Keyes, 1982). While the variations in elastic parameters are expected to be ~1% for the doping levels here, they can still have a significant effect on Rayleigh wave speed (c_r). To account for the change, LiSAW was performed on bare silicon wafers prior to film deposition. The results

were used to fit the c_{11} stiffness tensor for subsequent analysis in theoretical curve determination. The doped wafers were found to have $c_r = 4896 \text{ m/s}$ in the [100] direction which is a bit lower than the results for undoped wafers ($c_r = 4917 \text{ m/s}$). This corresponded to -1%decrease in c_{11} lowering it from 165.7 GPa to 163.8 GPa. This value was used for all theoretical calculations to follow. The other silicon constants were assumed to be their theoretical values of $c_{12} = 63.9 \text{ GPa}$, $c_{44} = 79.6 \text{ GPa}$, and $\rho = 2330 \text{ kg/m}^3$.

To make the V₂O₅ surface more reflective for optical detection methods, a thin layer of aluminum was deposited on top of the V₂O₅ by magnetron sputtering (AJA International, Scituate, MA) as a secondary top film. The elastic modulus of the sputtered aluminum films, a key parameter for theoretical fitting, was confirmed by performing LiSAW and nanoindentation on separate samples without V₂O₅. Both techniques measured a modulus of 70 GPa which is in agreement with literature (Serway et al., 2008). The other properties of aluminum used for theoretical dispersion curves were $\rho = 2700 \text{ kg/m}^3$ and $\nu = 0.33$.

The first V₂O₅ film tested with LiSAW received a 100 nm aluminum top layer. SAWs were generated in the [100] direction of the silicon substrate and detected after traveling various distances ranging between 10 mm and 15 mm from the irradiation site. SAW waveforms obtained at 10 mm and 15 mm are shown in Fig. 6. Dispersion of the signals is evident from the waveform elongation on the 15 mm signal, especially visible at low frequencies. Spectral analysis showed a frequency response between 15 and 60 MHz in both signals. The upper bound of the detected frequencies is quite low considering that the detection spot size of <10 µm is sufficient to identify frequencies well in excess of 100 MHz. However, the nanostructured roughness was a likely cause of high frequency attenuation. It should be noted that the penetration depth of SAWs is on the order of one wavelength therefore a 50 MHz wave propagating on silicon has an influence of almost 100 µm. The film region represents only a small fraction of this influence, yet significant detectable dispersion occurs even at low frequencies. Thus, while a large bandwidth is desirable, it is not always necessary, as is the case in this sample where there was more than adequate bandwidth for dispersion analysis. This



Fig. 6. SAW signals recorded on the first V₂O₅ sample at a distance of 10 mm and 15 mm from the source. To aid in shot to shot SAW alignment, the **x** timescale is offset by 2 μ s such that sample irradiation occurs at \sim 2000 ns.

reasoning is why LiSAW is an attractive and effective technique for very thin films of delicate nature.

Experimental dispersion curves were calculated for the signals in Fig. 6 and the elastic modulus was extracted through fitting of theoretical curves. Four film parameters are necessary for fitting: Poisson's ratio, film thickness, density, and elastic modulus. Since we have linear dispersion, three of the four parameters must be known in advance. Fortunately, the Poisson's ratio (0.3 was used) has a negligible influence on dispersion (Xiao and You, 2006) and the thickness $(1.7 \,\mu\text{m})$ can be measured to a high degree of certainty. Density can typically be determined by methodologies involving Archimedes principle such as the case in this study where a pycnometer was utilized. The measured density was 2.28 ± 0.11 g/cm³. As mentioned previously, the BJH pore volume was measured to be $0.054 \text{ cm}^3/\text{g}$ which suggests a film density of 2.84 g/cm^3 . While this is considerably higher than the value obtained by pycnometry, it assumed that all internal sample pores were intruded on in the powder sample. Any discrepancy would lower the density further. It should be noted that both measurements are considerably less dense than an ideally packed V₂O₅ structure which has a density of 3.36 g/cm³ thus confirming substantial film porosity. The value given by pycnometry was utilized for analysis since it provides the best estimate of the bulk macroscale density for which LiSAW wave propagation characteristics are derived.

With all of the substrate and film parameters accounted for, the experimental dispersion curve was fit to a modulus of 53 ± 4 GPa. The error bars are based on the uncertainty in density measurement since this property dominates over other sources of error. The error was resolved by creating theoretical dispersion curves with the bounds given by the density measurement at a constant modulus of 53 GPa. These produced upper and lower bounding dispersion curves for which the elastic modulus was then fitted to thereby producing upper and lower modulus bounds. Fundamentally, the resolution in separation distance, spectral bandwidth of the probe laser, and time resolution of the detection electronics also contributes to error, but all of these combined would produce <1 m/s of error in phase velocity for a 100 MHz signal. This is substantially smaller than the 10+ m/s in error that the density uncertainty provides. Furthermore, curve fitting error was low and did not contribute significantly to error since the experimental curves were very linear and thus correlated well with the fit ($R^2 = 0.968$).

A second V_2O_5 film, with identical specifications as the first film, was tested to demonstrate measurement repeatability. On this sample, a 200 nm aluminum reflective layer was sputtered prior to testing. The SAWs were detected at smaller propagation distances, 5–11 mm, in an effort to capture higher frequencies before they attenuated. This resulted in a slight increase to the frequency limit with a new range of 15–75 MHz. The experimental dispersion curves generated from combinations of signals detected at 5 mm, 10 mm, and 11 mm are shown in Fig. 7. Theoretical curves were generated and the best fit curve (Fig. 7) revealed a modulus of 52 GPa which was in excellent agreement to the first sample.

3.3. Mechanical properties by nanoindentation

Nanoindentation was performed on the first of the two LiSAW samples described above. The test preceded the LiSAW study and was conducted before the reflective aluminum layer was deposited. The in-situ scanning probe microscopy (SPM) capability of the indentation system allowed for surface analysis prior to indentation. An SPM topography image is shown in Fig. 8. The roughness of the film is visibly evident and the measured RMS roughness resulting from a 10×10 µm area was 40 nm. Manual polishing was attempted with 0.05 um alumina powder on a low nap fine polishing cloth (Ted Pella, Redding, CA), but the film delaminated and detached too easily. Thus while the film/substrate interface integrity is good enough to survive LiSAW, it could not survive the polish process even with just the sample holder weight of approximately 100 g applied. As a result, nanoindentation were performed on the as-synthesized state without polishing.



Fig. 7. Dispersion curves and theoretical fit for the second V₂O₅ sample.



Fig. 8. $10 \times 10 \,\mu\text{m}$ SPM image of the V₂O₅ surface: (a) 2D topography, (b) 3D profile.

Initially, nanoindentation proceeded to depths of below 80 nm in order to simulate characterization of the thickest film for which electrochemical data was reported on previously. This corresponds to \sim 10% (in order to minimize substrate effects) of the 800 nm thick film used for the electrochemical analysis in this study. For thin film electrodes, this 800 nm thickness is still quite large since most are considerably thinner (i.e. 3 min, 125 nm film) in order to minimize Li-ion diffusion time and thus maximize power capability (Wang and Cao, 2008). Indentations were performed ranging from 13 to 74 nm in maximum depth and the calculated film modulus was 50 ± 10 GPa. The error bars represent the 1σ value resulting from a total of 74 measurements across many areas on the sample surface. It is apparent that the roughness is a likely culprit for the high standard deviation since many of the indentations were comparable in size to surface asperities. Previous measurements on a smooth film with the same tip, similar indentation depths, and a similar film moduli produced a 1σ error of <1.5 GPa (Johnson et al., 2007).

It should be noted that the modulus extracted from indentation loading curves is known as the reduced modulus (or indentation modulus), E_r , and its value represents a combination of film and indenter tip properties. Eq. (4) can be used to isolate the film modulus, E_f , requiring the tip properties (modulus $E_i = 1140$ GPa and Poisson's ratio $v_i = 0.07$ for the diamond Berkovich tip) and the film Poisson's ratio (v_f), but unlike LiSAW, the extraction is quite sensitive to the Poisson's ratio.

$$\frac{1}{E_r} = \frac{(1 - v_i^2)}{E_i} + \frac{(1 - v_f^2)}{E_f}$$
(4)

A Poisson's ratio of 0.3 (identical to LiSAW fitting) was used to obtain the 50 GPa value reported above, but considering a range of v = 0.2 - 0.35, the modulus can differ by as much as 47.9–52.4 GPa. Thus, uncertainty in the Poisson's ratio can affect the nanoindentation measurement error in addition to the normal site to site error; in this case, the additional error was 4–5%. A similar argument can be made about the LiSAW technique which requires and is sensitive to the film density. However, density is typically less challenging to procure by direct means and the site to site error is lower for LiSAW since the measurement naturally averages over the SAW propagation distance.

A second round of indentation was performed to larger depths in order to explore whether the roughness could be overcome. The maximum depth was extended to 170 nm representing 10% of the film thickness. The indentations revealed a modulus of 51 ± 9 GPa, a slight increase compared to the previous tests with an improvement on the error. The increase is well within the error bars of the previous test, but could be explained by possible collapse of the porous network at shallow indentation depths prior to densification at larger depths (Lu et al., 2012). Different from our previous work on a thin porous zeolite film (Chow et al., 2014), no substrate effects were witnessed in the current study since the film was sufficiently thick. Overall, the nanoindentation data produced results similar to the LiSAW tests, but with substantially higher measurement error. Considering the film roughness of 40 nm, even the deeper indentations could easily be affected by the surface features. This highlights one of the advantages of the LiSAW technique by demonstrating accurate analysis on a rough nanostructured film.

4. Conclusion

From the studies performed, it is clear that both LiSAW and nanoindentation are valuable techniques in the field of thin film mechanical electrode materials characterization. They were both able to measure the elastic modulus of the porous nanostructured V₂O₅ electrode, but LiSAW was more accurate on the rough film. Two LiSAW samples were tested resulting in an elastic modulus of 53 ± 4 GPa for the first sample and 52 ± 4 GPa for the second. Both samples exhibited very linear dispersion curves which correlated well to the theoretical fit, thus the major contribution to error came from uncertainty in the film density. Nanoindentation produced a modulus of 50 ± 10 GPa at shallow indentation depths and 51 ± 9 GPa at depths up to 170 nm, or 10% of the film thickness. Roughness had a pronounced effect as evident by the large standard deviations. Uncertainty in the film Poisson's ratio, which is necessary to extract elastic modulus from indentation loading curves, can contribute an additional 4-5% error

for this film if v = 0.2 - 0.35 (v = 0.3 was used) are considered. The slightly lower results from nanoindentation could be the result of indentation induced pore collapse at low loading. However it is noted that the results are all within the error bars of one another, although the uncertainty of nanoindentation measurements was considerably larger. Overall, while the results extracted from both techniques were similar, LiSAW showed distinctive advantage as it did not encounter the challenges of the nanostructured roughness. For thinner films, the advantages of LiSAW are even greater since the technique naturally accounts for substrate influence.

The results of this work have large implications towards emerging thin film energy materials. As evolving knowledge in this field continues to drive films in the direction of more complex structures, compositions, and geometries, new methodologies for mechanical characterization will inevitably be required. The robustness and performance of materials are frequently linked to mechanical properties, thus the ability to characterize mechanical behavior is critical in development. LiSAW is a promising tool for this application since it has shown capability on complex films. It also has great utility and versatility since it can selectively target the more difficult properties such as density and Poisson's ratio. As shown in this paper, the intricacy of future energy materials will undoubtedly stress current mechanical characterization techniques and the spectrum of film tests will undoubtedly benefit from the proficiencies afforded by LiSAW. Specific to electrode materials, it is hoped that the technique can be applied to many more emerging electrode thin films such that a deeper understanding of the relationship between mechanical properties and electrode performance can be discovered.

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