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# Magnetically Detected Protein Binding Using Spin-Labeled Slow Off-Rate Modified Aptamers

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**ABSTRACT:** Recent developments in aptamer chemistry open up opportunities for new tools for protein biosensing. In this work, we present an approach to use immobilized slow off-rate modified aptamers (SOMAmers) site-specifically labeled with a nitroxide radical via azide—alkyne click chemistry as a means for detecting protein binding. Protein binding induces a change in rotational mobility of the spin label, which is detected via solution-state electron paramagnetic resonance (EPR) spectroscopy. We demonstrate the workflow and test the protocol using the SOMAmer SL5 and its protein target, platelet-derived growth factor B (PDGF-BB). In a complete site scan of the nitroxide over the SOMAmer, we determine the rotational mobility of the spin label in the absence and presence of target



protein. Several sites with sufficiently tight affinity and large rotational mobility change upon protein binding are identified. We then model a system where the spin-labeled SOMAmer assay is combined with fluorescence detection via diamond nitrogen–vacancy (NV) center relaxometry. The NV center spin–lattice relaxation time is modulated by the rotational mobility of a proximal spin label and thus responsive to SOMAmer–protein binding. The spin label-mediated assay provides a general approach for transducing protein binding events into magnetically detectable signals.

**KEYWORDS:** aptamer, protein assay, EPR spectroscopy, spin labeling, quantum sensing

egular and complete proteomics panels can provide real-Ntime diagnostic insight into human health and enable preventive and early treatments of nascent disease states.<sup>1,2</sup> Such measurements must be reliable and require fast, accurate, and accessible multiprotein sensors. Novel designs for such sensors featuring small sample volumes and high sensitivity have become feasible due to the recent developments in aptamer chemistry and quantum technologies.<sup>3-5</sup> A special class of DNA aptamers known as slow off-rate modified aptamers (SOMAmers) have been used as molecular binding reagents in a variety of biochemical studies.<sup>6-9</sup> SOMAmers are selected for tight binding interaction to target proteins and contain a series of nucleobases modified using hydrophobic constituents such as benzene and naphthalene rings.<sup>10</sup> The incorporation of these modifications has been demonstrated to improve the binding specificity and affinity of SOMAmers. With these properties, SOMAmers are reliable agents for capturing traces of specific target proteins.<sup>11,12</sup> In general, the SOMAmers are selected to form tight interactions with their targets with  $K_d$  values in the range of several nM to pM.<sup>13–16</sup> Structural information has been obtained from several SOMAmer-protein complexes.<sup>6,17-19</sup>

The specificity and affinity of SOMAmer-protein interactions can be fully utilized if they are paired with sensitive detection methods, ideally at the single-molecule level. One option for such a quantum sensor is the negatively charged nitrogen vacancy (NV) center in diamond. The NV center is a defect color center with a spin-triplet (S = 1) ground state localized on the dangling carbon bonds next to the vacancy. It can be spin polarized via optical pumping and provides a spinsensitive fluorescence readout. The use of the NV center as a quantum sensor for local magnetic fields is well studied.<sup>20–22</sup> Numerous studies have demonstrated the ability of NV centers to detect the presence of proximal spin species, such as proton spins on the diamond surface,<sup>23</sup> proton spins in the solvent,<sup>24–26</sup> and single electron spin centers in proximal biomolecules.<sup>27–32</sup>

There are several challenges that must be overcome in order to integrate SOMAmers and NV center-mediated detection into a device for protein sensing. One is to develop a biocompatible surface on NV center-hosting diamond for SOMAmer attachment.<sup>33</sup> Another is translating the SO-MAmer–protein binding event to a magnetic signal, which then can be detected via differential optical readout of an NV center. To address the second challenge, here we propose a protein binding assay that uses site-directed spin labeling of SOMAmers to generate a magnetic signal from a SOMAmer–

 Received:
 January 17, 2023

 Accepted:
 May 25, 2023

 Published:
 June 10, 2023







**Figure 1.** Schematic representation of NV-detected protein binding using spin-labeled SOMAmers. A. The apo and holo states of a SOMAmer covalently attached to the surface of a functionalized diamond via a linker (red). A spin label at distance *r* from the NV center is rotating with rate constant  $\tau_c$  with little restriction. Protein binding potentially restricts and slows the rotational mobility and changes the distance from the spin label to the NV center. B. Chemical structures of uracil (black) modified with dibenzocyclooctyne (DBCO, purple) and nitroxide spin label (green). The asterisks indicate rotatable bonds. The spin label is incorporated via catalyst-free azide–alkyne cycloaddition.

protein binding event. Figure 1A illustrates the principles of this proposed mechanism. The unbound SOMAmer (apo state) is spin-labeled with a paramagnetic nitroxide radical (S = 1/2) at a specific nucleotide and is covalently attached to the diamond surface, proximal to an NV center. The spin label is incorporated at an appropriately chosen solvent-exposed site of the SOMAmer and is therefore rotationally mobile. When a target protein binds to a SOMAmer (holo state), the spin label's rotational mobility is reduced, and its distance to the NV center is possibly shifted. Both the rotational mobility change and the distance change affect the NV center spin–lattice relaxation rate, characterized by the longitudinal relaxation time constant  $T_1^{.24,26,34}$  This rate can be measured using an optically detected magnetic resonance (ODMR) approach.<sup>24,35–37</sup>

For this detection mechanism to work, the spin-labeled SOMAmer must satisfy several prerequisites. First, the incorporation of the spin label must not degrade the SOMAmer–protein binding affinity; second, the spin label must undergo a significant rotational mobility or distance change upon protein binding; third, the rotational mobility or distance change must be on a scale that sensitively affects the NV center longitudinal relaxation rate. To identify sites that optimally satisfy these requirements, a complete scan of all spin labeling sites, i.e., all nucleotides, on a SOMAmer is necessary. While distance changes are only relevant in the presence of NV centers, rotational mobility changes can be measured using solution-state electron paramagnetic resonance (EPR) spectroscopy, a technique sensitive to molecular rotation rates on a time scale between 0.1 and 100 ns.<sup>38</sup>

In this work, we present a solution protein binding assay based on immobilized spin-labeled SOMAmers. We use a nitroxide-based spin label, since nitroxides are well-studied in solution EPR experiments and are sensitive reporters of molecular rotational motion.<sup>39,40</sup> The spin label is incorporated into a SOMAmer via site-specific azide-alkyne cycloaddition as shown in Figure 1B.<sup>41,42</sup> A model system was used for optimizing the spin labeling procedure and demonstrating the workflow for selecting labeling sites using two criteria: a small dissociation constant ( $K_d$ ) and a large change in rotational mobility in response to protein binding, quantified by the rotational correlation time constant ( $\tau_c$ ) of a fitted Brownian rotational diffusion model. The model system consists of a SOMAmer named SL5 and its target protein, platelet-derived growth factor B (PDGF-BB dimer).<sup>17</sup> We performed binding affinity assays for all labeling sites on SL5 and identified 19 nondisruptive sites. The associated constructs were immobilized on agarose beads via streptavidin—biotin linkers. The rotational mobility changes were then evaluated using solution-state EPR spectroscopy combined with spectral simulations to obtain  $\tau_c$  before and after addition of PDGF. Several sites show a significant change  $\tau_c$ , demonstrating the feasibility of this magnetically detected assay. Using the experimentally determined  $\tau_c$  we also model the impact of  $\tau_c$  on NV center  $T_1$  relaxation and discuss the effect of the NV center—spin label geometry on the overall robustness of a  $T_1$  relaxometry-based detection mechanism.

#### MATERIALS AND METHODS

Spin-Labeled SOMAmer Preparation. An ABI 3900 automated DNA synthesizer (Applied Biosystems) was used with conventional phosphoramidite methods with minor changes to the coupling conditions for modified phosphoramidites. Modified phosphoramidites were used in 0.1 M solutions using acetonitrile with 0-40% dichloromethane and 0-20% sulfolane as the solvent. Solid support was an ABI-style fritted column packed with controlled pore glass (CPG, LGC Biosearch Technologies) loaded with 3'-DMT-dT succinate with 1000 Å pore size. All syntheses were performed at the 50 nmol scale and the 5' end of each sequence was modified with a hexaethylene glycol (HEG) spacer and biotin group for support attachment. Introduction of a DBCO-modified nucleotide variant was done as a single-base replacement at selected sites within the DNA strand using phosphoramidites synthesized by SomaLogic, Inc. Deprotection was accomplished by treatment with concentrated ammonium hydroxide at 55 °C for 4-6 h, after which the product mixtures were filtered and residual solvents removed in a Genevac HT-12 evaporator. Identity and percent full-length product were determined using an Agilent 1290 Infinity LC system with an Agilent 6130B single quadrupole mass spectrometry detector using an Acquity C18 column 1.7  $\mu$ M 2.1 × 100 mm (Waters).

The resulting crude DBCO-modified SOMAmer residues were then redissolved in water for injection (WFI, HyPure WFI Quality Water, HyClone Laboratories, or equivalent) to 0.17 mM concentration (based on synthesis scale). A 100 mM solution of commercially sourced 4-azido-2,2,6,6-tetramethylpiperdinyl-1-oxyl (TEMPO) azide (Glen Research 50–2007–92) was prepared in dimethyl sulfoxide. Each oligonucleotide mixture received an aliquot of the azide solution at a 4:1 ratio of azide to SOMAmer (based on synthesis scale) and the resulting mixture was mixed at room temperature for 24 to 65 h, at which time analysis by LC/MS (Agilent 1290 Infinity, configured as above) confirmed that each cycloaddition reaction had reached completion (SI Figure S1). Each reaction mixture was then centrifugally filtered (Millipore Amicon Ultra-15



**Figure 2.** Evaluation of the binding affinity of spin-labeled SL5. A. A single SL5-PDGF-B complex from the crystal structure (PDB ID 4hqu).<sup>17</sup> B. Binding affinity assay of SL5 spin-labeled at all nucleotides. The bottom row indicates the nucleotide sequence in the unlabeled construct, where "M" represents any non-canonical nucleotide. The  $K_d$  of unlabeled construct is 6 pM, indicated on the left in gray. The typical error in  $\log_{10}(K_d)$  is about ±0.15, i.e., an overall uncertainty for  $K_d$  of a factor of 2. The green shaded region represents the acceptable range (under 6 nM) with opacity representing tightness of binding. Sites with low affinity (paler color bars) were excluded from EPR measurements.

3K), washed three times with 5 mL WFI per wash for removal of small molecule impurities. Product was collected in approximately 500 L WFI without further purification.

To spin label residues 1-4 and 11-29, we used the SL5 parent construct that contains a HEG linker in place of residues 5-10. To spin label residues 5-10, we used SL5 with the sequence GACZAC in place of the HEG linker, where Z is the modified nucleotide 5-benzylaminocarbonyl.<sup>17</sup> In this paper, the numbering is based on the 29-residue length.

**Binding Affinity Measurements.** To determine the  $K_d$  of modified aptamers, filter binding assays were performed using the SB18T binding buffer (40 mM HEPES, 102 mM NaCl, 5 mM KCl, 5 mM MgCl<sub>2</sub>, and 0.01% Tween-20, pH 7.4). Modified aptamers were 5' end labeled using T4 polynucleotide kinase (New England Biolabs) and  $\gamma$ -[<sup>32</sup>P]ATP (PerkinElmer). Radiolabeled aptamers (~20,000 cpm) were mixed with PDGF-BB (Creative BioMart) at concentrations ranging from 10<sup>-7</sup> to 10<sup>-12</sup> M and incubated at 37 °C for 40 min. Following incubation, 5.5 mL of 400 mg/mL silica-based microspheres (5 M Zorbax PSM 300A, Cat. No. 899999–555, Agilent Technologies) was added, and bound complexes were captured on Durapore filter plates (EMD Millipore). The fraction of bound aptamer was quantified with a phosphoimager (Typhoon, GE Healthcare), and data were analyzed in ImageQuant (GE Healthcare).

Sample Preparation for EPR Measurements. The SOMAmer constructs used in EPR measurements were biotinylated  $^{11}$  at the  $5^\prime$ end to facilitate immobilization on agarose beads coated with immobilized recombinant streptavidin protein (Pierce Streptavidin Plus UltraLink #53117; ThermoFisher Scientific). The bead diameter varies from 50 to 80  $\mu$ m. SOMAmer concentration was determined by measuring the absorbance at 260 nm. Resin (0.5 mL of 50% slurry) was buffer exchanged into SB18T buffer, the supernatant was removed, and 20 nmol of SOMAmer was applied to the resin. The reaction was incubated, with end-over-end rotation, at 37 °C for 1 h. Resin was pelleted and supernatant was removed and discarded. Unbound streptavidin was quenched with 0.9 mL reaction buffer containing biotin. SOMAmer-loaded resin was washed thrice with 1 mL reaction buffer and resuspended in a final volume of 0.5 mL for a final concentration of 40 µM SOMAmer in 50% slurry. SOMAmerloaded resin was stored at 4 °C. The PDGF-BB protein was stored at -80 °C and then dissolved before use.

**EPR Measurements.** Continuous-wave (CW) EPR experiments were performed with 40  $\mu$ M SOMAmers in SB18T buffer. For the protein binding experiments, protein targets with concentration of 45–50  $\mu$ M were added. For each sample, 12–15  $\mu$ L were transferred into quartz capillaries (1.0 mm O.D. and 0.7 mm I.D.; Sutter Instrument). Solution-state EPR spectra were collected at room temperature (23 °C) using an X-band continuous-wave (CW) Bruker EMX spectrometer equipped with a ER 4123D dielectric resonator with a resonance frequency of 9.77 GHz. All spectra were obtained with a peak-to-peak modulation amplitude of 0.1 mT and a sweep rate of 0.24 mT/s and were collected under nonsaturating conditions at

0.64 mW. Power saturation data are provided in the Supporting Information (SI Figure S2).

**EPR Spectral Simulations.** All spectra were aligned to the same magnetic field range and scaled by number of scans, receiver gain, and square root of power. Then, background was removed by subtracting the spectrum of buffer solution only. The spectra were integrated and normalized to obtain absorption spectra for simulation and comparison. An isotropic rigid-body Brownian rotational diffusion model<sup>43,44</sup> without motional restriction was fitted to the measured spectral lineshapes. The isotropic rotational correlation time  $\tau_c$  and the *g*-tensor principal values ( $g_{xy} g_{yy} g_z$ ) were varied for each fit. The hyperfine tensor principal values ( $[A_{xx} A_{yy} A_z] = [21.79(3), 12.43(2), 103.01(2)]$  MHz) and the homogeneous broadening were simulated from a spectrum collected using a frozen solution of a free SOMAmer labeled at site 1 (SI Figure S4). Spectral simulations were performed using the open-source MATLAB toolbox EasySpin 6.0.0-dev.49,<sup>45</sup> with an example MATLAB script given in the SI (SI Section 2.3).

#### RESULTS

Binding Affinity Assay. To identify nondisruptive labeling sites, the binding affinities for PDGF of all SL5 constructs spinlabeled at each of the 29 sites were measured. The determined values of  $K_d$  are shown in Figure 2. The measured  $K_d$  of the unlabeled SOMAmer is 6 pM (indicated on the left in Figure 2). No spin-labeled construct binds to PDGF as tightly as the unlabeled parent SOMAmer. Based on a previously obtained crystal structure of the SL5-PDGF complex, the labeling sites structurally fall into three domains: a direct binding domain, an indirect binding domain, and a nonbinding domain.<sup>1</sup> Incorporation of a spin label into any site in the nonbinding domain (yellow) does not significantly disrupt the interaction since all  $K_d$  values in this region are under 100 pM. For the direct and indirect binding domains (purple and blue, respectively), no clear correlation between the physical location and the magnitude of  $K_d$  degradation is observed. Some sites are highly disruptive (e.g., 18 and 23), whereas others (12, 13, 16) have  $K_d$  values similar to those with labeled sites in the nonbinding domain. For the purpose of selecting optimal labeling sites for NV center detection, any construct with a  $K_d$  larger than 6 nM (1000-fold reduction compared to parent construct) was deemed too disruptive and excluded from further analysis. Site 29 was excluded as well, as spin labeling at this site was unsuccessful.

**EPR Measurements and Rotational Mobility Analysis.** To magnetically detect protein binding with the spin-labeled SOMAmers in solution, room-temperature EPR spectra were recorded before and after protein binding (SI Figure S3



**Figure 3.** Evaluation of the mobility change of spin-labeled SL5 upon protein binding. A. An example of the integrated spectra for the apo (blue) and holo (orange) states of SL5 labeled at site 12 with simulated lineshapes obtained by fitting a rotational diffusion model. Fitted  $\tau_c$  values are given in corresponding colors. The black arrow marks the feature that is most indicative of the mobility change. B. Control experiments using site 12 with nontarget proteins IgG (red) and transferrin (yellow). C. Fitted  $\tau_c$  values of for the apo and holo states of all EPR-evaluated constructs and for two nontarget protein controls. The excluded sites are indicated by shaded columns with colors corresponding to Figure 2.

bottom). The change in spectral shape is significant and is due to a change in the overall molecular size. This demonstrates the feasibility of this EPR-based solution binding assay. To mimic a diamond surface for potential optical readout of this assay, 19 biotinylated constructs that passed the binding affinity screening were ligated to streptavidin-coated agarose beads. This suppresses the global tumbling of the entire SOMAmer or SOMAmer-protein complex (SI Figure S3 top) and isolates the rotational motion of the spin label. Again, spectra were recorded in the absence and presence of target protein (SI Figures S5 and S6). All spectra were analyzed as described above, and rotational correlation times  $au_{\rm c}$  were determined for all sites (SI Tables S2). Figure 3A shows an example using the construct labeled at site 12. Spin label rotational diffusion slows down after addition of the target protein PDGF as indicated by the broadening of the spectral line shape in the low-field region around 346 mT (indicated by an arrow). This region of the spectrum corresponds to the  $m_{\rm I}$  = +1 manifold of the <sup>14</sup>N nucleus of the nitroxide. Additionally, the high-field edge of the spectrum at about 351 mT sharpens slightly. Least-squares fits of simulated spectra to the experimental data reveal a significant increase in  $\tau_{c}$ . The results for all sites, shown in Figure 3C, indicate that significant changes in  $\tau_c$  are observed for many sites.

To evaluate the possible change in  $\tau_c$  from nonspecific bindings, control experiments were performed using the SOMAmer construct labeled at site 12 and two different nontarget proteins: immunoglobulin G (IgG) and transferrin. The two nontarget proteins were chosen for their high abundance in human plasma.<sup>46</sup> The results shown in Figure 3B indicate that the addition of nontarget proteins does not significantly change  $\tau_c$ . Small shifts are observed at the maxima of the holo spectra of PDGF and nontarget proteins, possibly indicating changes in the *g*-tensor of the electron spin. The shifts could be attributed to the changes in the hydrogenbonding environment around the nitroxide group of the spin label.<sup>47</sup> The slight increase in fitted  $\tau_c$  from the apo to the holo state of the control experiments may be a result of the simple model used in the fitting.

To compare both binding affinity and mobility changes observed from all sites, Figure 4 plots the values of  $\tau_c$  before and after protein binding for all constructs along with their  $K_d$ as measured by the binding affinity assays. The experimentally



**Figure 4.** Evaluation of the performance of all spin-labeled SOMAmer constructs based on  $K_d$  (vertical axis) and  $\tau_c$  (horizontal axis). The magnitude of the difference in  $\tau_c$  between the apo (×) and holo ( $\blacktriangleright$ ) states is indicated by arrow length, and the corresponding  $K_d$  values are indicated by arrow color according to the color bar on the left.

determined  $\tau_c$  are indicated by the corresponding markers, and the length of the arrow represents the magnitude of change in  $\tau_c$  (denoted as  $\Delta \tau_c$ ) after addition of protein. For the experiments using the target protein PDGF, sites 1, 2, 3, and 16 have the smallest  $\Delta \tau_c$ , whereas sites 5, 7, 9, and 11 have the largest. Site 24 is the only one that shows a decrease of rotational mobility upon binding; nonetheless, the difference is relatively small. Optimal labeling sites should have a large  $\Delta \tau_c$ and a small  $K_d$ . Sites 5, 7, 9, and 12 best fit these criteria. Similar to the binding affinity screening results, there is no obvious correlation between the magnitude of rotational mobility change and the physical location of the labeling site, since sites 5, 7, and 9 are all in the nonbinding domain and sites 12 and 13 are in the direct binding domain.

Overall, the results indicate that spin labels can be used as reporters of SOMAmer-protein binding and can therefore form the basis of protein sensors with high sensitivity and specificity.



**Figure 5.** Predicted dependence of NV center  $T_1$  on measured  $\tau_c$ . A. Simulating the  $T_1$  of the NV center as a function of  $\tau_c$ . Top: estimated  $T_1$  of an NV center 6 nm from the spin label, with a fixed angle  $\theta = 45^{\circ}$ . Blue solid line: the minimum  $T_1$  at  $\tau_c = 0.055$  ns. Gray dotted line: the intrinsic  $T_1$  of the NV center in the absence of a spin label,  $T_1^0$ . Bottom: the spectral density evaluated at  $\omega_{NV}$ . Pink solid line: the maximum of the spectral density at  $\tau_c = 1/\omega_{NV}$ . B. The change in spectral density at  $\omega_{NV}$  as a function of and apo and holo  $\tau_c$ . The color bar indicates the differences between density amplitudes at two  $\tau_c$  values. Yellow dots:  $\tau_c$  values of apo and holo states of SL5 constructs. Pink solid line indicates  $1/\omega_{NV}$ . C.  $T_1$  (ms) of the NV center a function of NV center—spin label distance and orientation. The contour lines represent NV center  $T_1$  in the presence of a mobile spin label with  $T_1^0 = 3.5$  ms and  $\tau_c = 4$  ns at a location specified by distance r and angle  $\theta$ .

#### DISCUSSION

Next, we explore how the EPR-detected SOMAmer protein binding assay demonstrated above can be combined with fluorescence detection via diamond NV centers. Having experimentally determined the rotational correlation times  $\tau_c$ , we can estimate the effect of protein binding on the  $T_1$ relaxation time of an NV center in proximity to the spinlabeled SOMAmer using a previously published model.<sup>24,26</sup> In this model, the total  $T_1$  relaxation rate of an NV center in proximity to a spin label is a sum of three independent contributions:

$$\frac{1}{T_1} = \frac{1}{T_1^0} + \frac{1}{T_1^{\text{rot}}} + \frac{1}{T_1^{\text{rlx}}}$$
(1)

Here,  $1/T_1^0$  is the NV center's intrinsic spin–lattice relaxation rate in the absence of the spin label,  $1/T_1^{\text{rot}}$  is the contribution from the rotational dynamics of the spin label, and  $1/T_1^{\text{rbx}}$  is the contribution from the spin label's own intrinsic spin relaxation dynamics.

The first term depends on temperature,  $^{48,49}$  diamond surface quality,  $^{50-52}$  depth of NV center from surface,  $^{37,53}$  and the spin environment,  $^{54}$  especially the  $^{13}C$  concentration in the diamond lattice.  $^{55,56}$  For a single NV center implanted in a  $^{12}C$  isotopically purified diamond,  $T_1^0$  can be as long as several milliseconds  $^{37,57,58}$  at room temperature, whereas bulk NV centers in nanodiamonds may have  $T_1^0$  as short as hundreds of  $\mu$ s at room temperature.  $^{26,49,59}$ 

The second term represents the effect of a fluctuating magnetic field at the NV center due to the rotational diffusion dynamics of the spin label, mediated by the dipolar coupling between spin label and NV center. At the appropriate frequency, this fluctuating field induces transitions at the NV center. The associated transition rate is given by

$$\frac{1}{T_1^{\text{rot}}} = 3\gamma_{\text{NV}}^2 B_{\perp,\text{rot}}^2 J_{\tau_c}(\omega_{\text{NV}})$$
(2)

where  $\gamma_{\rm NV}$  is the gyromagnetic ratio of the NV center,  $B_{\perp,\rm rot}$  is the averaged magnetic field generated by the spin label in the plane perpendicular to the magnetization direction of the NV center, and  $J_{r_c}(\omega_{\rm NV})$  is the spectral density of the rotating spin label evaluated at  $\omega_{\rm NV}$ , the angular frequency for the  $m_{\rm s} = 0 \leftrightarrow$ ±1 transition of the NV center. The full derivation of this relation is given in the Supporting Information (SI section 3.2).  $B^2_{1,rot}$  is given by

$$B_{\rm L,rot}^{2} = \frac{1}{4} \left( \frac{\mu_{0}}{4\pi} \hbar \gamma_{\rm SL} \right)^{2} \frac{2 + 3 \sin^{2} \theta}{r^{6}}$$
(3)

where  $\mu_0$  is the vacuum magnetic permeability,  $\gamma_{SL}$  is the gyromagnetic ratio of the spin label, *r* is the distance between the spin label and the NV center, and  $\theta$  is the angle between the magnetization direction of the NV center in a weak external magnetic field and the NV–spin label direction. Via *r* and  $\theta$ ,  $B^2_{\perp,rot}$  contains the geometry dependence of  $1/T_1$ : the rate drops off rapidly with increasing distance, but is only mildly orientation dependent as the  $\theta$  term varies beween 2 for  $\theta = 0$  and 5 for  $\theta = \pi/2$ .

The spectral density in eq 2 is given by

$$J_{\tau_{\rm c}}(\omega) = \frac{\tau_{\rm c}}{1 + \tau_{\rm c}^2 \omega^2} \tag{4}$$

and represents the fluctuating field amplitude as a function of frequency  $\omega$ . Its value at  $\omega = \omega_{NV}$  is used in eq 2.

The third term in eq 1 is the contribution from the relaxation dynamics of the spin label, which is relayed to the NV center by the same mechanisms as the second term. It contains several separate contributions:<sup>60,61</sup>

$$\frac{1}{T_1^{\text{rlx}}} = \left(\frac{\mu_0}{4\pi}\right)^2 \gamma_{\text{NV}}^2 \gamma_{\text{SL}}^2 \hbar^2 \frac{1}{r^6} \left(\frac{1}{6}B + 3C + \frac{3}{2}E\right)$$
(5)

$$B = \frac{I_{2,SL}}{1 + (\omega_{NV} - \omega_{SL})^2 T_{2,SL}^2} (1 - 3\cos^2\theta)^2$$
(6)

$$C = \frac{T_{\rm l,SL}}{1 + \omega_{\rm NV}^2 T_{\rm l,SL}^2} \sin^2 \theta \, \cos^2 \theta \tag{7}$$

$$E = \frac{T_{2,SL}}{1 + (\omega_{NV} + \omega_{SL})^2 T_{2,SL}^2} \sin^4 \theta$$
(8)

where  $\omega_{SL}$  is the (angular) transition frequency of the spin label electron, and  $T_{1,SL}$  and  $T_{2,SL}$  are the spin–lattice and spin–spin relaxation times of the spin label, respectively. Each of the three terms *B*, *C*, and *E* is a product of a spectral density (depending on  $T_{1,SL}$  or  $T_{2,SL}$ ) and an angular term. Given a reasonable NV-spin label geometry (r = 6 nm and  $\theta = 45^{\circ}$ ) and typical ranges of relaxation times of nitroxides, 0.3 to 3 s for  $T_1$  and 0.1 to 1 s for  $T_{2r}^{62,63}$  values for  $1/T_1^{\text{rlx}}$  are significantly smaller than the  $T_1^0$  relaxation rate (details are discussed in SI section 3.3). Therefore, the  $1/T_1^{\text{rlx}}$  term is negligible and is dropped from now on.

Figure 5 elaborates on the dynamic and geometric aspects of  $1/T_1^{\text{rot}}$ . Figure 5A illustrates the effect of  $\tau_c$  for a fixed NV-spin label geometry (r and  $\theta$ ). The top panel shows estimated NV  $T_1$  as a function of spin label  $\tau_c$  from 0.01 to 50 ns, using the first two terms in eq 1 with  $T_1^0 = 3.5$  ms, r = 6 nm, and  $\theta = 45^\circ$ . The bottom panel shows the spectral density from eq 4 as a function of  $\tau_c$ . The spectral density can be used to infer the magnitude of  $T_1$  reduction. The density amplitude quantifies the transition rate at the NV center due spin label reorienting at different rates. To isolate the dynamics of the spin label from the geometric factors, Figure 5B illustrates the effect of the  $\tau_{c}$ values for the apo and the holo states on the change in spectral density (which is proportional to the change in relaxation rate). A large change in spectral density will enhance  $T_1$ contrast. The largest spectral density changes occur if either  $au_{
m c}^{
m apo}$  or  $au_{
m c}^{
m holo}$  is close to  $1/\omega_{
m NV}$  and the other is at least an order of magnitude slower or faster (white and black regions in Figure 5B). The apo/holo  $\tau_c$  values for SL5–PDGF are shown by the cluster of yellow dots. The  $\tau_c$  values of apo and holo state range 2.5-5 ns and 6-16 ns, respectively. They do not fall within the most sensitive range. One approach to improve upon this is to build more flexible linkers into the structure by exploring a wider range of spin label structures and coupling chemistries.

Figure 5C visualizes the dependence of  $T_1$  on the relative position of the NV center spin and the nitroxide spin for a fixed  $\tau_c$ . According to eq 3,  $T_1^{\text{rot}}$  contains an angular dependence and a factor of  $1/r^6$ , and the relaxation enhancement is most pronounced if the spin-spin distance is short and the angle  $\theta$  perpendicular. The contour plot shows  $T_1$  using  $T_1^0 = 3.5$  ms and  $\tau_c = 4$  ns. The reduction in  $T_1$ increases as the angle  $\theta$  approaches 90° and decreases with increasing distance. This is an important aspect for combining the spin-labeled SOMAmer assay with a diamond sensor since both orientation and distance affect the sensitivity of the  $T_1$ relaxation-based sensing mechanism.

To evaluate the reduction of  $T_1$  relaxation rate due to the protein binding-induced changes in rotational mobility for spin-labeled SL5,  $T_1$  values are estimated for all sites using the same geometry (r = 6 nm and  $\theta = 45^{\circ}$ ) as used in Figure 5A. The calculated results are ranked from the largest to the smallest binding-induced  $T_1$  change from left to right in Figure 6. In this geometry, the construct labeled at site 11 produces the largest contrast in  $T_1$ , followed by sites 21 and 10. However, sites 5, 7, 9, and 12, which yield the largest differences in  $\tau_c$ , are not among the highest-ranked sites in terms of  $T_1$  contrast. This is because the apo state  $\tau_c$  of these constructs are longer than those of sites 11, 21, and 10 (Figure 4), resulting in a shorter window of  $\Delta \tau_c$  as illustrated by Figure 5A.

The estimated  $\Delta T_1$  ranges from 0.05 ms (site 11) to less than 0.01 ms (site 1 through 4). Recent work has demonstrated a reliable resolution of 100  $\mu$ s using nanodiamonds.<sup>64</sup> Several publications also managed to distinguish  $\Delta T_1$  values of hundreds of ms to tens of  $\mu$ s.<sup>59,65–67</sup> The SL5-PDGF binding-induced reduction of the estimated NV center  $\Delta T_1$  falls in a range that can be detected by current NV center



**Figure 6.** Evaluating the estimated  $T_{1,NV}$  for all tested sites, assuming r = 6 nm and  $\theta = 45^\circ$ . The magnitude of the difference in  $T_{1,NV}$  between the apo (×) and holo ( $\blacktriangle$ ) states is indicated by arrow length, and the corresponding  $K_d$  is indicated by arrow color with a reference in Figure 4.

sensor devices.<sup>24,26,32,36,68,69</sup> Nevertheless, several strategies can be applied to increase  $T_1$  contrast and therefore the overall sensitivity of the NV center relaxometry.

NV center  $T_1$  relaxometry detection sensitivity is strongly limited by the NV-spin label distance due to the steep  $1/r^6$ dependence in eq 3. Therefore, keeping r as short as possible is crucial. The depth of the NV center below the diamond surface can be engineered via tuning the energy of ion implantation.<sup>70,71</sup> The caveat is that shallow NV centers are often susceptible to magnetic noise from surface spins.<sup>51,52,57</sup> Another approach to potentially increase  $\Delta T_1$  is to shorten the spin label linker length by using alternative labeling chemistries such as coupling of azide-functionalized spin labels with terminal alkynes<sup>72</sup> or strained cycloalkynes,<sup>73</sup> or by ligation to the phosphate group,<sup>74</sup> the sugar,<sup>75</sup> and the base.<sup>76–79</sup> A series of rigid spin labels that can be covalently conjugated to the base were designed for measuring the flexibility of DNA;<sup>76,80-84</sup> however, these labels are less useful in this context, as their rigidity prevents the detection of spin label mobility changes. The  $\tau_c$  screening protocol reported in this work relies on the built-in mobility from the rotatable bonds in the linker (Figure 1B). In addition to linker length and rigidity, linker size is another critical aspect. An inflexible and bulky linker may have greater potential to disrupt binding interactions in comparison to a small and compact structure. The ideal linker would be short but flexible so that the spin label is close to the SOMAmer and responsive to protein binding but can still span a wide dynamic range of rotational motion.

The affinity of SOMAmer-protein binding may affect the overall sensitivity of the NV center detection. This effect is difficult to isolate and must be evaluated with the changes in spin label mobility upon protein binding. However, the SOMAmers may undergo conformational changes when binding to target proteins, potentially leading to distance changes. According to eq 3, the effect from a distance change can be as dramatic as the effect from a change in  $\tau_c$ . The distance change may counteract or enhance the effect of the  $\tau_c$  change on  $T_1$ . Further investigations of SOMAmer structure and conformational landscapes, also in the context of their attachment to a diamond surface,<sup>33</sup> will provide more information on probable distance effect.

# CONCLUSIONS

We developed a novel protocol for sensing proteins using immobilized site-specifically spin-labeled SOMAmers in combination with CW EPR spectroscopy to detect changes in spin label rotational mobility upon protein binding. A full site scan on the benchmark SOMAmer SL5 and its target protein PDGF-BB revealed a number of suitable sites where (a) the binding-induced change in rotational mobility is significant and (b) the reduction in binding affinity due to the presence of the spin label is small. The approach presented here is general and can be extended to other SOMAmer–protein pairs.

In addition, we modeled the detection of the observed rotational mobility changes via  $T_1$  relaxometry of a proximal diamond NV center, which provides a platform for increased sensitivity and multiplexing. The results indicate several strategies for maximizing contrast in NV  $T_1$ . First,  $T_1$  is minimized in the unbound state by shortening the rotational correlation time to approach the inverse angular transition frequency of the NV center. Second, the slow-down of  $\tau_c$  experienced by the spin label upon protein binding is maximized. Both strategies can be pursued by varying the nature of the spin label and scanning across labeling sites. Also, the NV-spin label distance needs to be carefully controlled in order to maximize  $T_1$  contrast.

# ASSOCIATED CONTENT

#### **1** Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acssensors.3c00112.

Details of spin-labeled SOMAmer characterization, EPR experiments and derivation of the expression for  $T_1$  of the NV center (PDF)

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# Notes

The authors declare no competing financial interest.

# ACKNOWLEDGMENTS

This work was supported by the National Science Foundation (ITE-2040520 to L.G., J.P.C. and S.S.; CHE-2154302 to S.S.).

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